# Crystal structures of fluorinated fullerenes

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 $C_{60}F_x$  and  $C_{70}F_x$ , respectively, prepared by the direct reactions of  $C_{60}$  and  $C_{70}$  powders with elemental fluorine form face-centered-cubic lattices with  $a_0=17.19$  and 17.93 Å, in which rotational and/or orientational disorders of the  $C_{60}F_x$  and  $C_{70}F_x$  molecules prevail.

### **1. INTRODUCTION**

Recent structural studies on solid  $C_{60}$  have established that it forms a face-centered-cubic (fcc) structure at room temperature and that the molecules execute rotational diffusion [1,2]. Structural studies of crystalline  $C_{70}$  have indicated that the equilibrium phase above room temperature is fcc, but that structural defects as well as residual solvent tend to stabilize the hexagonal-closest-packed (hcp) phase [3]. We have reported that fluorinated  $C_{60}$ ,  $C_{60}F_x$ (20<x<45), in which the fluorine atoms are bound to the carbon atoms, form an fcc lattice with  $a_0 \approx 17.1$  Å [4].

In this paper the crystal structure of  $C_{70}F_x$  is reported in connection with the crystal structure of  $C_{70}$ , and comparisons were made with the lattice parameters of  $C_{60}$  and  $C_{60}F_x$ .

## 2. EXPERIMENTAL

 $C_{60}$  and  $C_{70}$  powder samples, heat-treated for 2 days in a dynamic vacuum, were reacted with elemental fluorine of 1 atm at 200° C for 4 hours and at RT for 4 weeks, respectively. The compositions of the samples were determined by chemical analysis on F employing the oxygen flask combustion method with the fluoride ion electrode. The Debye-Scherrer (DS) X-ray powder diffraction (XRPD) patterns were obtained by loading powder samples in a quartz capillary ( $\phi$ =0.7 mm) using Cu-K $\alpha$  radiation on the Shimadzu X-ray diffractometer XD-D1. The intensities were collected in  $0.1^{\circ}$  (2 $\theta$ ) steps in the range 5-35° (2 $\theta$ ) using a pinhole slit ( $\phi$ =2 mm).

#### **3. RESULTS AND DISCUSSION**

The compositions of the products were ca.  $C_{60}F_{45}$  and  $C_{70}F_{32}$ . The DS-XRPD patterns for  $C_{60}$ ,  $C_{60}F_{45}$ ,  $C_{70}$  and  $C_{70}F_{32}$  are shown in Figures. 1(a)-4(a) and their lattice constants  $a_0$  (and  $c_0$ ) are summarized in Table 1.

The simulated pattern (Figure. 1(b)) for  $C_{60}$  was obtained assuming rotational and/or orientational disorders of  $C_{60}$  molecules, *i.e.*, spherical shell (SS) model [2,3] with the center to carbon-atom distance of 3.52 Å [2]. For the simulation the Gaussian function with an appropriate standard deviation was applied to

Table 1 Lattice constants for  $C_{60}$ ,  $C_{60}F_{45}$ ,  $C_{70}$  and  $C_{70}F_{32}$ 

| Sample                          | Lattice    | $a_0/\text{\AA}$ | <i>c</i> ₀/Å |
|---------------------------------|------------|------------------|--------------|
| C <sub>60</sub>                 | fcc        | 14.19            |              |
| $C_{60}F_{45}$                  | fcc        | 17.19            |              |
| C <sub>70</sub>                 | hcp<br>fcc | 10.53<br>14.90   | 17.21        |
| C <sub>70</sub> F <sub>32</sub> | fcc        | 17.93            |              |

each calculated peak intensity.

The simulated pattern (Figure. 2(*b*)) for  $C_{60}F_{45}$  was obtained assuming a single unique center to carbon-atom distance of  $3.87(=3.52\times1.1)$  Å and a single unique center to fluorine-atom distance of  $5.23(=3.52\times1.1+1.36)$  Å, where the skeleton radius was expanded by the factor of 1.1, and the C-F distance of 1.36 Å was employed [5,6]. The sphere expands owing to the  $\pi$ -electron depletion. The atomic scattering factors of C

**Figure 1.** Observed (a) and simulated (b) diffraction patterns for  $C_{60}$ .

20

 $2\theta / \deg$ 

10

(a)

(b)

30



**Figure 2.** Observed (*a*), simulated (*b*) and calculated (*c*) diffraction patterns for  $C_{60}F_{45}$ .

and F were weighted by 60 and 45, respectively, to match the composition. The molecule consists of two concentric shells: one made of carbon atoms and the other of fluorine. It is noted that any  $C_{60}F_x$  molecule except  $C_{60}F_{60}$  will have several non-equivalent carbon atoms, and some non-equivalent fluorine atoms depending on the molecular symmetry. As a consequence a rotating  $C_{60}F_x$  becomes essentially equivalent



**Figure 3.** Observed (*a*) and simulated (*b*) diffraction patterns for  $C_{70}$ , and calculated patterns for hcp (*c*) and fcc (*d*).



**Figure 4.** Observed (a), simulated (b) and calculated (c) diffraction patterns for  $C_{70}F_{32}$ .

ntensity (arb. units)

to multi-shells of carbon and fluorine atoms. However, the satisfactory fit in Figure. 2 indicates that the differences in their radii are not significant.

The structure of  $C_{70}$  is more subtle. Recent studies have indicated that heat-treatment of  $C_{70}$ is not sufficient to remove the trapped solvent completely [3]; sublimation and the subsequent annealing of  $C_{70}$  is necessary to obtain an fcc crystal of pure  $C_{70}$ . Our  $C_{70}$  sample was found to be an admixture of hcp and fcc phases, indicating that there remained a certain amount of solvent in the  $C_{70}$  sample. Figure 3 shows the calculated patterns based on the hcp (c) and fcc (d) structures and the simulated pattern (b)with a 2:1 mixture of the hcp and fcc phases. The calculation was done assuming the SS model. There are five unique center to carbonatom distances; 4.172 Å (10 atoms), 4.029 Å (10), 3.876 Å (20) 3.663 Å (20) and 3.565 Å (10) [7].

The XRPD (Figure. 4(a)) of  $C_{70}F_{32}$ , on the other hand, indicates that it forms an fcc lattice. The structural change from hcp to fcc for  $C_{70}$  upon fluorination is intriguing. The pattern was simulated (Figure. 4(b)) in the same manner as above; the  $C_{70}$  skeleton was enlarged by the ratio 1.1 and 70 fluorine atoms were attached 'normal' to the spheroid. The C–F bond length of 1.36 Å was used. The atomic scattering factors of C and F were weighted by 70 and 32, respectively. The molecule consists of ten concentric shells: five of carbon atoms and five

**Table 2** Effective radii of spherical  $C_{60}$ ,  $C_{60}F_{45}$ ,  $C_{70}$  and  $C_{70}F_{32}$ , and the increments upon fluorination

| System               | <br>Radius/Å | Increment/Å |  |
|----------------------|--------------|-------------|--|
| C <sub>60</sub> 5.02 |              | 1.06        |  |
| $C_{60}F_{45}$       | 6.08         | 1.00        |  |
| C <sub>70</sub>      | 5.27         |             |  |
| $C_{70}F_{32}$       | 6.34         | 1.07        |  |

of fluorine.

In Table 2 the effective radii of spherical  $C_{60}$ ,  $C_{60}F_{45}$ ,  $C_{70}$  and  $C_{70}F_{32}$  are summarized. For  $C_{70}$  the hexagonal  $c_0/a_0$  ratio is 1.63, suggesting an ideal hcp structure made of perfect spheres. The same differences of *ca.* 1.07 Å in the radii of the pristine fullerenes and the fluorinated ones indicate that the fluorinations of  $C_{60}$  and  $C_{70}$  bring about similar changes in the molecular as well as in the crystal structures of the fullerenes.

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