# Investigation of Sc<sup>3+</sup> in LiNbO<sub>3</sub> single crystals

Joyce K. Yamamoto and K. Yamagishi

Mitsui Mining and Smelting Co. Ltd., Corporate R and D Center 1333-2 Haraichi, Ageo, Saitama 362 Japan

The effect of trivalent  $Sc^{3+}$  doping in LiNbO<sub>3</sub> was studied and compared to  $Mg^{2+}$ -doped LiNbO<sub>3</sub>. The 1 to 3 mol% Sc-doped LiNbO<sub>3</sub> single crystals were grown by the Czochralski technique. The effect of scandium dopant concentration on the lattice parameters, Curie temperature and laser (515 nm) induced photorefractive damage was investigated.

### 1. INTRODUCTION

Since its discovery LiNbO<sub>3</sub> has been a well known and well studied ferroelectric material. It is characterized by a large spontaneous polarization,  $P_S = 0.71 \text{ C} \cdot \text{m}^{-2}$ , which gives rise to large linear and nonlinear optical coefficients. This has led to many applications in piezoelectric, surface acoustic wave, acoustooptic and electro-optic devices.

One area that LiNbO<sub>3</sub> could have an impact would be as a frequency doubler for infrared radiation, utilizing the large nonlinear dcoefficients,  $d_{31} / ld_{36}^{KDP}I = -12.3$  and  $d_{31} / ld_{36}^{KDP}I = -8.6$ , to efficiently produce green light. [1]

This application is hampered by the easy susceptibility of LiNbO<sub>3</sub> to changes in its refractive indices. When laser radiation is incident on the crystal, the electro-optic coefficients couple to space charge-generated electric field. This is known as photorefractive or optical damage and it is especially dominant in the more energetic visible wavelength range.

Due to the high carrier conductivity in the illuminated region, the charge carriers accumulate at the periphery of this region forming, forming an electric field,  $E_j$ . From the equation:

$$\Delta \left(\frac{1}{n^2}\right)_i = \sum_j r_{ij} E_j \tag{1}$$

the induced changes in the refractive indices are determined by the linear electro-optic rcoefficients. In LiNbO<sub>3</sub> the  $r_{33}$  (32 x 10<sup>-12</sup>m\*V<sup>-1</sup>) is 3 times larger than the  $r_{31}$  (10 x 10<sup>-12</sup>m\*V<sup>-1</sup>) coefficient, therefore the observed birefringence,  $\Delta n = n_e - n_o$ , is primarily due to changes in  $n_e$ . [1]

#### 1.1. Mg-DOPED LITHIUM NIOBATE

It has been reported that MgO doping up to 4.5 to 5.0 mol% drastically reduces the photorefractive damage. [2] The reduction in the damage level is reportedly two orders of magnitude below the undoped LiNbO<sub>3</sub>.

Subsequent studies determined that in this concentration region there is a change in the defect structure which in turn influences the photoconductivity. [3] Other physical properties were also found to change at this MgO concentration. [4] At low MgO concentrations, the Curie temperature,  $T_{c'}$ , increased, peaking in the 5-7 mol% range and again decreased in the high concentration range. This behavior has been attributed to a change in Mg site occupancy from the A-site to

the B-site, resulting in a maximum vacancy concentration corresponding to the peak in the Tc.

There has been a vast amount of work into dopant effects in LiNbO<sub>2</sub>, where the largest interest has been in photorefractive applications, where a maximum change in birefringence is needed to efficiently produce holographic images in the crystal. This is possible with multi-valent transition metal cations, where the photon induced oxidation/ reduction process serves as a source of excited state carriers/trapping sites. Since it is evident that the defect structure and dopant play an important role in determining the response to laser radiation, the effect a trivalent cation with one valence state has on LiNbO3 may provide a basis for comparison to the divalent dopants to provide a better understanding of the dopant-defect-property relationship.

## 1.2. Sc-DOPED LITHIUM NIOBATE

Scandium was selected because the cationic radius is is very similar to  $Mg^{2+}$  and from the distribution coefficient, the affinity of the LiNbO<sub>3</sub> structure for Sc and Mg are about equal, Table 1. Work into scandium doped LiNbO<sub>3</sub> was first reported in 1966, where the as-grown domain structure was studied. [5] Later the Li<sub>2</sub>O-Nb<sub>2</sub>O<sub>5</sub>-Sc<sub>2</sub>O<sub>3</sub> phase diagram was investigated. [6] Recently, scandium was found to reduce the photorefractive damage in LiNbO<sub>3</sub>, leading to studies into the concentration effects on the optical and physical properties. [7]

## 2. EXPERIMENTAL

The Sc-doped  $LiNbO_3$  crystals were grown by the Czochralski technique and pulled from congruent (Li/Nb = 0.941) melts doped with 1

Table 1. LiNbO<sub>3</sub> Dopant Cations

Cation	Valence	lonic Radii* (nm)	Keff**
Li	1+	7.4	
Nb	5+	6.4	
Mg	2+	7.2	1.3
Sc	3+	7.3	1.2
Ti	4+	6.05	0.7-0.9

\* From Ref. 11

\*\* From Ref. 1

to 3 mol% Sc added in their oxide forms. The crystals, 2.5 cm in diameter and 5 - 6 cm in length, were inclusion-free, water-clear and transparent. The measured distribution coefficient of 1.3 was close to the reported value of 1.2. [1] Samples were taken from subgrain boundary-free area of the crystal, as determined by X-ray topography. Growth striations and inhomogeneities were not observed in both Schlieren and polarized microscopy techniques.

The lattice parameters were calculated from a least squares fit of 32 high two-theta peaks, using a silicon internal standard. The Curie temperature, was measured using differential thermal analysis. The experimental procedure was the same as described in the literature. [8] Chemical analysis was performed using X-ray fluorescence to measure the scandium concentration and inductively coupled plasma technique was used to measure the iron concentration, which was less than 1 ppm in all Powdered samples for lattice crystals. parameter and Curie temperature measurements were taken from the same area in the crystal as that used for compositional analysis.

### 3. RESULTS AND DISCUSSION

The lattice parameters as a function of scandium concentration were found to be in good agreement with the literature and generally followed the same trends as Mg:LiNbO<sub>3</sub>, Figure 1. The effect of the increasing dopant concentration was to smoothly and gradually increased the size of the unit cell, with a predominant basal plane expansion. As previously pointed out in the literature, this result is not unexpected since the cation sizes of the dopant and the host are approximately the same.

The  $T_c$  of the Sc:LiNbO<sub>3</sub> exhibited the same general behavior as Mg:LiNbO<sub>3</sub>, but the magnitude of the  $T_c$  elevation was smaller, and the peak in the  $T_c$  occurred at a much lower Sc concentration, Figure 2. Due to the similarity of the two dopants in terms of size, the  $T_c$ elevation may also indicate Sc occupation on the A-site in the low scandium region and joint occupation of the A and B-sites in the high scandium region. Previously reported ESR results indicated that Fe<sup>3+</sup> occupies the A-site and therefore based on size and valence, it is reasonable to assume that the Sc cation would occupy the A-site in low concentrations. [10]

The photorefractive damage caused by 515 nm radiation from an argon ion laser was measured using the Senarmont compensator method, using a HeNe (633 nm) laser as the probe beam. In Figure 3, the 1 mol% Sc:LiNbO3 displayed the lowest level of damage. The damage level increased as the scandium concentration increased above the 1 mol% level. This is a departure from the behavior observed in Mg:LiNbO3. The minimal damage level in MgO:LiNbO3 reportedly occurs near 5 mol% Mg. A comparable damage level was observed to occur in Sc:LiNbO3 at the 2 mol% Sc level.

The cause of this behavior is not clearly understood. It has been proposed that an



Figure 1. Lattice parameters as a function of Sc concentration. Mg:LiNbO<sub>3</sub> data was taken from Ref. 12.

increase in vacancy concentration and decrease in photorefractive damage are related. The vacancies may serve as trapping sites for mobile carriers and by increasing the trapping site concentration the carrier mobility is decreased, resulting in a weaker electric field. [2] If it is assumed that the maximum in the  $T_c$  is due to a maximum in the vacancy concentration, which has been shown to occur, then the observed behavior in Mg:LiNbO3 does seems to agree with the vacancy model. [4] This also seems to be the case for the Sc:LiNbO<sub>3</sub>, where the maximum  $T_c$  does occur at 1.2 mol% Sc and the minimum damage level occurred at 1 mol% Sc. The difference in the dopant concentration indicates that the Sc cation is more efficient in producing cation vacancies than the Mg cation. The higher valence could create two vacancies for each Sc cation if the Sc occupies the A-site. Whether



Figure 2. Curie temperature as a function of Sc concentration



Figure 3. Photorefractive damage induced by argon ion laser (515 nm, 7 W/cm) of several Sc:LiNbO<sub>3</sub> single crystals. 633 nm HeNe laser was used as a probe beam.

the lower damage level is a direct consequence of the high vacancy concentration is still not known. More work into trivalent and divalent dopants should shed more light on the relationship between the vacancy concentration and photoconductivity.

#### 4. CONCLUSIONS

 $Sc_2O_3$  and MgO doping of LiNbO<sub>3</sub> was found to have qualitatively similar physical effects on the lattice parameters,  $T_c$  and photorefractive damage level as a function of dopant concentration. The effect of dopant valence was attributed to differences observed in the concentration dependence of the measured properties.

### REFERENCES

- 1. E. Kaldis (ed.), Current Topics in Materials Science, vol. 1, North-Holland, 1978.
- D. A. Bryan, R. Gerson and H. E. Tomaschke, Appl. Phys. Lett., 44 (1984) 847.
- K. L. Sweeney, H. E. Halliburton, D. A. Bryan, R. R. Rice, R. Gerson and H. E. Tomaschke, Appl. Phys. Lett., 45 (1984) 805.
- B. C. Grabmaier, W. Wersing and W. Koestler, J. Cryst. Growth, 110 (1991) 339.
- 5. K. Nassau, H. J. Levinstein and G. M. Loiacono, J. Phys. Chem. Sol., 27 (1966) 983.
- A. Baumgarte and R. Blachnik, J. Them. Anal., 35 (1989) 617.
- J. K. Yamamoto, K. Kitamura, N. Iyi, S. Kimura, Y. Furukawa and M. Sato, Appl. Phys. Lett., 61 (1992) 2156.
- P. K. Gallagher and H. M. O'Bryan Jr., J. Am. Ceram. Soc., 68 (1985) 147.
- P.K. Gallagher and H.M. O'Bryan Jr., J. Am. Ceram. Soc., 71 (1988) C56.
- 10. A. Boker, H. Donnerberg, O. F. Schirmer and F. Xiqi, J. Phys., C2 (1990) 6865.
- 11. R. D. Shannon and C. T. Prewitt, Acta Cryst., B25 (1969) 925.
- 12. Y. Furukawa, M. Sato, F. Nitanda and K. Itoh, J. Cryst. Growth, 99 (1990) 996.