## Electronic Structures of Strongly Correlated Itinerant-Electron System

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Fermi surfaces and mass enhancement factor in strongly correlated itinerant-electron compounds, CeCo<sub>2</sub>, CeRh<sub>2</sub>, CeRu<sub>2</sub>, YCo<sub>2</sub> and LuCo<sub>2</sub>, are calculated. The calculated Fermi surfaces of CeCo<sub>2</sub> and CeRu<sub>2</sub> are in good agreement with the dHvA measurements. Mass enhancement factors,  $\tilde{\gamma}$ , are calculated by using a method of perturbation expansion with respect to electron-electron interactions based on Fermi liquid theory. The values of electronic specific-heat coefficient calculated by using  $\tilde{\gamma}$  reproduce well the experimental values for CeCo<sub>2</sub>, CeRu<sub>2</sub>, YCo<sub>2</sub> and LuCo<sub>2</sub>.

Key word: mass enhancement factor, Fermi surface, heavy fermion compounds, cubic Laves phase

## 1 Introduction

Density functional theory (DFT) within the local density approximation (LDA) have provided good descriptions for electronic states in many compounds. Calculated Fermi surfaces in strongly correlated electron (SCE) system, delectron and / or f-electron compounds, are in good agreement with experimental ones [1]. It is naturally accepted because DFT describes correctly physical quantities including many-body effects in ground states such as Fermi surfaces. On the other hand, the values of physical quantities in low-energy excited states, such as an electronic specific-heat coefficient  $(\gamma)$  and cyclotron effective masses, are hardly reproduced [1]. The reason is that DFT do not guarantee the many-body effects in low-energy excitation. In SCE system, most of such effects are caused by electron-electron interactions between SCEs. It is important how to approach many-body effects in understanding the physical property in such system.

One of the method for considering the manybody effects has been given by Yamada *et al.* [2] In the method, a calculation of enhancement for  $\gamma$  is equivalent to obtain an enhancement factor for SCE,  $\tilde{\gamma}_k$ . A simple method to calculate  $\tilde{\gamma}_k$  is a second-order perturbation. However it is still difficult to calculate  $\tilde{\gamma}_k$  numerically in a threedimensional system, because it is necessary to carry out multiple integrals with respect to momentum. An enhancement factor can be calculated numerically only from a density of states (DOS) for SCEs if  $\tilde{\gamma}_k$  is replaced by  $\tilde{\gamma}(=\sum \tilde{\gamma}_k)$ , in which *k*-dependence is neglected. Series of the cubic Laves-phase compounds have shown various physical properties such as mass enhancement, superconductivity and metamagnetism. Band structure calcualtions have been reported for CeCo<sub>2</sub>, CeRh<sub>2</sub> and CeRu<sub>2</sub> [3]. Their Ce-f electrons are shown itinerant properties because of hybridization between Ce-f and conduction electrons, though be localized at the rare earth atom in general. The fermi surfaces of CeCo<sub>2</sub> and CeRu<sub>2</sub> are in good agreement with the dHvA experiments.

 $YCo_2$  and  $LuCo_2$  are well-known as a strongly enhanced Pauli paramagnets like Pd metal, in which Co-d electrons are strongly correlated ones. Band structure calcualtions have been reported for  $YCo_2$  and  $LuCo_2$  [4]. Recently, different characteristics of magnetic properties of  $CeCo_2$  and a group of  $YCo_2$  and  $LuCo_2$  have been experimentally revealed. <sup>59</sup>Co NMR measurement for  $YCo_2$ , [5, 6] for  $LuCo_2$  [7] and NMR and NQR measurements for CeCo<sub>2</sub> [8, 9] indicate: (1) the measured value of  $\chi$  associated with Co-3d for  $CeCo_2$  is much smaller than that for YCo<sub>2</sub> and LuCo<sub>2</sub>; (2)  $(T_1T)^{-1}$  of CeCo<sub>2</sub> is proportional to  $K_{3d}^2$ , where  $K_{3d}$  is a Knight shift with Co-3d electrons, but proportional to  $K_{3d}$  in  $YCo_2$  and  $LuCo_2$ . While the experimental values of  $\gamma$  ( $\gamma_{exp}$ ) are very similar in YCo<sub>2</sub>, LuCo<sub>2</sub> and CeCo<sub>2</sub>, the calculated values of  $\gamma$  ( $\gamma$ <sub>band</sub>) are much smaller than those. Then large enhancement in effective masses at the Fermi energy are expected.

This paper is reported the Fermi surfaces and the enhanced values of  $\gamma$  ( $\gamma_{ehs}$ ) for cubic Lavesphase compounds CeCo<sub>2</sub>, CeRh<sub>2</sub>, CeRu<sub>2</sub>, YCo<sub>2</sub>, LuCo<sub>2</sub>.

# 2 Method of Band Calculation

We have carried out band structure calculations by using a full-potential linear augmented plane wave (FLAPW) method. The crystal structure of compounds is a MgCu<sub>2</sub>-type which is called a cubic Laves-phase. The space group is  $Fd\bar{3}m$ . The exchange-correlation interactions of electrons are taken into account by the LDA, in the form proposed by Moruzzi et al. [10] We adopted the relaxed-core approximation to the Xe-core states except 5p electrons on Ce atom, the Ar-core on Co atoms and the Kr-core on Rh and Ru atoms. The scalar relativistic effects are taken into account for all electrons and the spinorbit interactions are included for valence electrons as a second variational procedure. The final potential is constructed self-consistently from eigenstates at 19 sampling k-points in the irreducible 1/48 Brillouin zone. To obtain a final band structure, we selected 231 sampling k-points in the irreducible 1/48 Brillouin zone.

### 3 Fermi Surfaces

We have calculated the Fermi surfaces with the extremal cross-sectional areas and the cyclotron effective masses for the compounds. The results are shown in our previous papers in detail [3]. The calculated Fermi surfaces for  $CeCo_2$ and  $CeRu_2$  can explain many observed branches and show similar topology. In  $CeRh_2$ , the uncertainty of assignment remains in several branches. It seems that the LDA calculation for  $CeRh_2$  fails to obtain Fermi surfaces measured by the dHvA experiment. The calculated cyclotron masses in heavy fermion compounds are, as usual, smaller than the observed ones, because of mass enhancement.

## 4 Density of States and Enhancement Factor

A theory for mass enhancement in strongly correlated system is given by Yamada *et al* [2] in the framework of Fermi liquid theory on the basis of a periodic Anderson model. The theory is a reasonable one to describe the essential properties of the system and can be applied for real compounds if the non-correlation term in Hamiltonian is replaced with that in DFT.

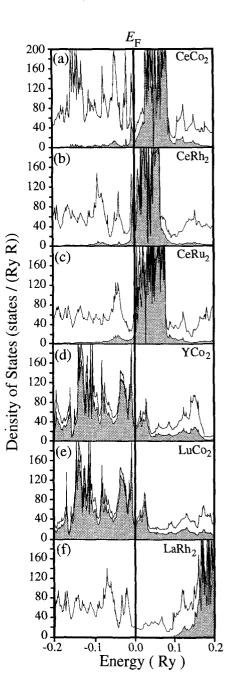


Figure 1: Calculated density of states (DOS) for (a)  $CeCo_2$ , (b)  $CeRh_2$ , (c)  $CeRu_2$ , (d)  $YCo_2$ , (e)  $LuCo_2$  and (f)  $LaRh_2$ . The solid lines and shadow areas show the total and Ce-f (for f electroncase) or Co-d partial DOS (for d electron-case), respectively.

Table 1: Results of DOSs and enhancement calculations with experimental data.  $\gamma_{\text{band}}$  is the electronic specific-heat coefficient, i.e., the total DOS ( $\rho(E_{\text{F}}^{o})$ ) obtained from a band calculation. R is the contribution of the component of SCE on Ce- or Co-site to the  $\gamma_{\text{band}}$ .  $\tilde{\gamma}$  is the enhancement factor according to partial DOS of SCE. (see text).  $\gamma_{\text{ehs}}$  is the value of  $\gamma$  taking into account the effects of  $\tilde{\gamma}$  ( $\gamma_{\text{ehs}} = \gamma_{\text{band}} \{1 + R_d(\tilde{\gamma} - 1)\}$ ).  $\gamma_{\text{exp}}$  is the experimental value.

	$\gamma_{ ext{band}} \ ( ext{mJ/K}^2 \cdot  ext{mol})$	$R \ (\%)$	$ ilde{\gamma}^\dagger$	${\gamma_{ m ehs} \over ({ m mJ}/{ m K}^2 \cdot { m mol})}$	$({ m mJ/K^2 \cdot mol})$
		f	electron case		
$\overline{\mathrm{CeCo}_2}$	14.6	33	4.76	33.3	35 [11]
$\operatorname{CeRh}_2$	11.9	60	7.60	59.4	$20 \ [12]$
$CeRu_2$	11.6	31	5.77	28.5	30 [13]
$CeSn_3$	11.7	61	7.90	60.7	53[1]
$LaRh_2$	3.7	3	1.08	3.7	3.7  [14]
		d	electron case		
YCo <sub>2</sub>	6.1	84	6.87	36.2	36.2 [15]
LuCo <sub>2</sub>	5.7	82	6.58	31.9	26.6 [16]
$CeCo_2$	14.3	43	1.38	$15.5 \ (33.3^{\dagger\dagger})$	35.0 [11]

†: calculated in the case of  $U_{ff}=5$  eV or  $U_{dd}=1.8$  eV

 $\dagger$ <sup>†</sup>: a value when  $U_{ff}$  between Ce-f electrons is considered

Following the theory, a coefficient of the T-linear term of the specific heat is

$$\gamma = \frac{\pi^2}{3} k_{\rm B}^2 \sum_k \left[ \tilde{\gamma}_k(\mu) \rho_k^l(\mu) + \sum_{\sigma} \rho_{k\sigma}^c(\mu) \right], \quad (1)$$

where

$$\tilde{\gamma}_k(\omega) = 1 - \frac{\partial \Sigma_k(\omega)}{\partial \omega},$$
(2)

 $\Sigma_k(\omega)$ , the self energy for the localized electrons. It is calculated by using the method of secondorder perturbation with respect to U. Unfortunately, it is difficult to calculate  $\partial \Sigma_k(\omega)/\partial \omega$  numerically for real compounds even in the secondorder perturbation. The reason is that it is necessary to carry out multiple (six-fold in a threedimensional system) integrals with respect to momentum k' and q.

On the other hand, a factor  $\tilde{\gamma}(=\Sigma\tilde{\gamma}_k)$  can be calculated easily compared to  $\tilde{\gamma}_k$ , because multiple integrals are replaced by three onedimensional integrals over an energy. We can use  $\tilde{\gamma}$ , when summations of  $\tilde{\gamma}_k$  and  $\rho_k^f(\omega)$  with respect to  $\boldsymbol{k}$  in eq.(1) are separated. Values of  $\gamma$  evaluated by  $\tilde{\gamma}$  are the same as those using  $\tilde{\gamma}_k$ , when  $\boldsymbol{k}$ -dependence of self-energy is neglected. Details for formulation of  $\tilde{\gamma}$  are shown in our previous paper [3]. The DOSs are shown in Figs. 1(a)- 1(f), respectively.

Using the partial DOSs and eq.(A.1) in ref. 3, , the enhancement factors  $\tilde{\gamma}$  are calculated. The value of U for f-electron is used 5eV, because the value is evaluated U~4-5eV from XPS and BIS experiments [17, 18] The results of the calculations are listed in Table 1. The main results are as follows:

(1) the values of  $\gamma_{\rm ehs}$  for Ce-compounds are larger than  $\gamma_{\rm band}$ , while the values of  $\gamma_{\rm band}$  and  $\gamma_{\rm ehs}$  for LaRh<sub>2</sub> is equal to  $\gamma_{\rm exp}$ ,

(2) the values of  $\gamma_{ehs}$  are fairly in good agreement with  $\gamma_{exp}$  except for CeRh<sub>2</sub>.

### 5 Summary

We have calculated band structures for CeCo<sub>2</sub>, CeRh<sub>2</sub>, CeRu<sub>2</sub> YCo<sub>2</sub> and LuCo<sub>2</sub> using the FLAPW method with spin-orbit interactions. Using the electronic structures, we have calculated DOS and Fermi surfaces for all the compounds. The calculated Fermi surfaces of CeCo<sub>2</sub> and CeRu<sub>2</sub> explain well the dHvA measurements. Using the calculated partial DOS for SCE and the method of a second-order perturbation with respect to many-body effects, we have calculated the values of  $\tilde{\gamma}$  and the enhanced electronic specific-heat coefficients  $\gamma_{ehs}$  for the above compounds. The results indicate that the values of  $\gamma_{\rm ehs}$  for CeCo<sub>2</sub>, CeRu<sub>2</sub>, YCo<sub>2</sub> and LuCo<sub>2</sub> reproduce the values of  $\gamma_{exp}$ . On the other hand, the calculated Fermi surfaces and the value of  $\gamma_{\rm ehs}$ for CeRh<sub>2</sub> show poor agreement with the experiments. Consequently, when we calculate the values of  $\gamma_{\rm ehs}$ , it is important to evaluate Fermi surfaces corresponding to the experiments.

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