γ '-Fe₄N Thin Films with Out-of-plane Magnetization

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Highly oriented γ' -Fe₄N(001) films on MgO(001) and SrTiO₃(001) substrates and γ' -Fe₄N(111) films on MgO(111) and SrTiO₃(111) substrates were successfully grown by means of reactive sputtering. Due to the structural and magnetic measurements, γ' -Fe₄N(111) films show an out-of-plane magnetic anisotropy and the direction of magnetization in the films depends on their substrates.

Key words: γ '-Fe₄N, thin film, reactive sputtering, conversion electron Mössbauer spectroscopy, magnetic anisotropy

1. INTRODUCTION

We are aiming to develop toxic-free magnetic devices. From such point of view, iron nitrides are very attractive material. They are made up of very common elements and both elements in themselves are not harmful to the environment. The nitrides have large corrosion and wear resistance, and moreover, some of these have excellent magnetic properties. Therefore, potential application for high-density magnetic storage devices is expected [1].

 γ '-Fe₄N is the most stable iron nitride phase, and is a ferromagnet which has large saturation magnetization (190emu/g) and the Curie temperature of 785K [2]. It adopts the cubic perovskite structure (a=3.795Å) with N atom located at the body-center of an fcc-Fe lattice. A Mössbauer spectrum of γ '-Fe₄N consists of three sextets which indicate ferromagnetic Fe atoms in one edge site and two magetically different face-center sites [3]. The easy axis of magnetization for γ '-Fe₄N is along <100> in both bulk and thin-film forms [3, 4].

Recently, epitaxial multilayers have attracted much attention with respect to magnetic devices like spin valve [5] and ferromagnetic tunnel junction [6]. In this study, we fabricated iron nitride films to obtain highly oriented γ' -Fe₄N for application to such devices, and investigated structural and magnetic properties of the films.

2. EXPERIMENTAL

Iron nitride thin films were fabricated by reactive RF magnetron sputtering using N₂. The base pressure in the deposition chamber was approximately 5×10^{-6} Pa. A RF magnetron sputter cathode and a N₂ gas inlet are located at angles of 60° and 30° to the surface of a substrate mounted on a rotating sample stage, respectively. The distance between the sputtering target and the substrate was 150mm.

We investigated hyperfine structure of γ' -Fe₄N films by conversion electron Mössbauer spectroscopy (CEMS). To acquire high signal/noise ratio for CEMS, ⁵⁷Fe enriched chips (95.7%, 3 pieces, about 6.5mm×6.5mm×0.6mmt for each piece) were placed on a pure Fe disk target (99.9%, 55mm ϕ ×2mmt) when γ' -Fe₄N films were deposited.

Samples were fabricated in the following steps A, B and C. In steps A and B, we sought an optimum

deposition condition for a γ' -Fe₄N single phase on silica glass by changing N₂ flow rates and substrate temperatures. Only the Fe disk was used as a target in step A, and the dummy chips were added in step B. Finally in step C, the ⁵⁷Fe enriched chips were placed on the disk instead of the dummy chips and ⁵⁷Fe enriched γ' -Fe₄N films were deposited on silica glass, MgO(001), MgO(111), SrTiO₃(001) and SrTiO₃(111) substrates by the optimum condition derived in step B. The pressure and flow rate of Ar gas were kept at 0.2Pa and 10.0sccm, respectively. The RF power applied to the target was 100W. The thickness of the films was about 2300Å.

Crystalline structure of the films was investigated by X-ray diffraction (XRD) with Cu-K α radiation. Surface morphology of the films was observed by scanning electron microscopy (SEM). Magnetic properties of the films were investigated by measuring the magnetic moment versus magnetic field (M-H) hysteresis curves by vibrating sample magnetometer (VSM). The hyperfine structure of the films was examined by CEMS with 57 Co(Rh) source. The velocity was calibrated using α -Fe. All of these measurements were performed at room temperature.

3. RESULTS AND DISCUSSION

3.1 Growth condition

The XRD patterns of iron nitrides deposited on silica glass in step A and B are shown in Fig.1. In step A, γ' -Fc₄N single phase is successfully grown at the temperature of 600K and the flow rate of 1.0sccm (Fig. 1(a)). However, α -Fe is included in the film fabricated in step B at the optimum condition obtained in step A (Fig. 1(b-1)). The increase in surface area of the target led to the increase in sputter rate. So the N₂ flow rates are set at 1.1sccm and the substrate temperatures are gradually changed from 600K to 800K. A γ' -Fe₄N film with small, constant amount of ε -Fe₃N was fabricated up to 700K (Fig.1(b-2)). The ε phase was decreased at 750K and disappeared at 800K (Fig.1(b-3)).

The peak intensity of γ '-Fe₄N films fabricated in step A and B is different from each other. The strong (200) and (111) peaks are seen in Fig.1(a), while the relative peak intensity of each peak in Fig.1(b-3) becomes similar to those of bulk samples. This is due to thermal fluctuation.

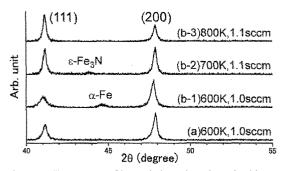


Fig.1: XRD patterns of iron nitride films deposited in step A and B.

The films are also observed by means of SEM. Both films consist of crystallites that are about 100 nm in diameter. The crystallites grown in step A have sharp edges and are close-packed. By contrast, those grown in step B are round and have gaps between each other. With such morphology, the γ' -Fe₄N film deposited in step B is easier to be oxidized.

3.2 Structure and surface morphology of γ '-Fe₄N films

Fig.2 shows the XRD patterns and SEM images of γ' -Fe₄N films grown on silica glass, MgO(001) and MgO(111) at the optimum condition found in step B (γ' /glass, γ' /MgO(001) and γ' /MgO(111), respectively). γ' /MgO(001) and γ' /MgO(111) are highly oriented and have the same orientation as their substrates. Both films, especially γ' /MgO(001), have very large grains. The crystallites as can be seen in γ' /glass and straight cracks are contained in γ' /MgO(111), but are not contained in γ' /MgO(001). Highly oriented γ' -Fe₄N films are also grown on SrTiO₃(001) and SrTiO₃(111) at the same condition (γ' /STO(001) and γ' /STO(111), respectively).

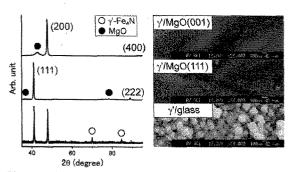


Fig.2: XRD patterns and SEM images of $\gamma^{\prime}\mbox{-}Fe_4N$ films on glass and MgO substrates.

Table I: Lattice parameters of $\gamma'\text{-}Fe_4N$ films, bulk $\gamma'\text{-}Fe_4N$ and MgO.

	d ₀₀₁ (Å)	d ₁₁₁ (Å)	d ₁₀₀ (Å)	α (deg.)
γ'/glass	3.804	2.198	3.791	89.82
γ'/MgO	3.815	2.201	3.785	89.74
γ'/STO	3.811	2.198	3.787	89.80
bulk γ'-Fe ₄ N	3.795	2.191	-	-
MgO	4.211	2.432	-	-
SrTiO ₃	3.905	2.253	-	-

Table II: Difference in interplanar spacings between γ '-Fe₄N films on single crystals, bulk γ '-Fe₄N and substrates.

substrates.					
Substrate		MgO		SrTiO ₃	
Form of γ'-Fe ₄ N		film	bulk	film	bulk
d (Å)	d ₀₀₁	3.815	3.795	3.811	3.795
	d ₁₁₁	2.201	2.191	2.198	2.191
d _{film} /d _{bulk}	d ₀₀₁	1.005		1.004	_
	d ₁₁₁	1.005		1.003	_
$d_{substrate}/d_{\gamma'}$	d ₀₀₁	1.104	1.110	1.025	1.029
	d ₁₁₁	1.105	1.110	1.025	1.027

The lattice parameters of γ' -Fe₄N films are shown in Table I. The out-of-plane interplanar spacings of the films d₀₀₁ and d₁₁₁ are derived from XRD patterns. The in-plane interplanar spacing d₁₀₀ for tetragonally packed (001) planes and the angle between edges of an unit cell α for rhombohedrally packed (111) planes are calculated on the assumption that the volume and the length of an edge of a unit cell are the same as bulk, respectively. All the films are nearly cubic. However, their out-of-plane interplanar spacings d₀₀₁ and d₁₁₁ are larger than those of bulk, and γ' /MgO and γ' /STO have larger d₀₀₁ and d₁₁₁ than γ' /glass.

The difference in interplanar spacings between the films on single crystals, bulk γ' -Fe₄N and substrates are listed in Table II. The ratio between interplanar spacings of γ' -Fe₄N and the substrates are 1.110 and 1.029, that is to say, 11 and 41 γ' -Fe₄N unit cells correspond to 10 MgO and 40 SrTiO₃ unit cells, respectively. The 11 and 41 bulk γ' -Fe₄N unit cells are slightly smaller than corresponding substrate unit cells. The lattice of the films are expanded about 0.3-0.5% than bulk. This lattice expansion makes the ratio between interplanar spacings of γ' -Fe₄N and the substrates lower. Thus, the lattice mismatch between γ' -Fe₄N films and their substrates is remarkably reduced by the distortion in γ' -Fe₄N films.

3.3 Magnetic properties of y'-Fe₄N films

The in-plane and out-of-plane M-H curves of γ' -Fe₄N films on glass and MgO substrates are shown in Fig.3. The values of saturation magnetization M_S and coercive force H_C for each film are listed in Table III. The effects of demagnetization field are not taken into consideration. All the films have smaller M_S than that of bulk.

Table III: Magnetic properties of y'-Fe4N films.

	in-plane		out-of-plane		
	M _S	H _C	M _S	H _C	
	(emu/g)	(Oe)	(emu/g)	(Oe)	
γ'/glass	151	92	128	147	
γ'/MgO(001)	166	35	129	. 77	
γ ['] /MgO(111)	160	46	154	42	

Two γ'/MgO films are sharply magnetized at low field when in-plane magnetic field is applied. At that time, $\gamma'/MgO(001)$ has the smallest H_c and the largest M_s. The values of in-plane and out-of-plane H_c in $\gamma'/MgO(111)$ are almost equivalent. $\gamma'/glass$ has the largest H_c in both in-plane and out-of-plane M-H curves.

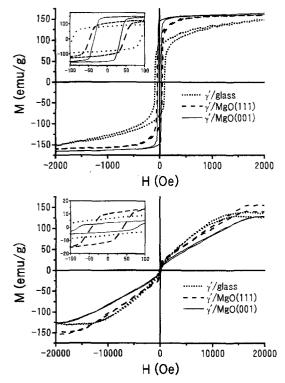


Fig.3: In-plane and out-of-plane M-H curves of y'-Fe₄N films on glass and MgO substrates. The insets show the area around zero field.

These results are due to the morphology of the films. Spins are easily arranged in highly oriented films with large grains like in the cases of γ'/MgO . The behavior of magnetization of γ'/MgO depends on their crystalline orientation. In contrast, the randomly oriented crystallites in y'/glass make domain walls hard to move.

3.4 Hyperfine structure of γ' -Fe₄N films

The CEMS spectra and hyperfine parameters of $\gamma'\text{-}Fe_4N$ films are shown in Fig.4 and Table IV, respectively. Each spectrum is composed of three sextets

Table IV: Hyperfine parameters of γ' -Fe₄N films.

		HF	QS	IS	RA
		(kOe)	(mm/s)	(mm/s)	(%)
γ'/glass	FeI	339.07	0.01	0.12	25
	FeIIA	215.85	0.03	0.25	50
	FeIIB	216.03	- 0.07	0.05	25
w/M~O	FeI	336.81	0.00	0.12	25
γ'/MgO (001)	FeIIA	216.33	0.19	0.19	50
(001)	FeIIB	214.60	- 0.38	0.18	25
γ'/MgO	FeI	334.34	0.01	0.11	25
	FeIIA	214.34	0.10	0.20	50
(111)	FeIIB	214.29	- 0.20	0.16	25
WETO	FeI	336.30	0.00	0.12	25
γ'/STO (001)	FeIIA	215.70	0.21	0.19	50
(001)	FeIIB	214.87	- 0.42	0.18	25
γ'/STO (111)	FeI	335.14	0.01	0.12	25
	FeIIA	211.45	0.16	0.21	50
	FeIIB	209.97	- 0.32	0.14	25

which indicate ferromagnetic Fe sites in γ' -Fe₄N.

The intensity of the peaks around ±2mm/s varies according to the substrates. In contrast, no substrate dependence is seen in hyperfine parameters. It is well known that the intensity ratio of three sub-doublets constituting one sextet depends on θ_{R} , the angle between sample magnetization and incident γ -ray [7]. The relationship between θ_S , the angle of sample magnetization to the substrate surface ($\theta_s = 90^{\circ} \cdot \theta_R$), and intensity of sub-doublets $I_{(i,j)}$ are given by:

$$I_{(1,6)}: I_{(2,5)}: I_{(3,4)} = 3: \frac{4\cos^2\theta_s}{1+\sin^2\theta_s}: 1.$$

The sample magnetization is normal to the substrate surface when $I_{(2,5)}/I_{(3,4)}$ equals to zero, and if $I_{(2,5)}/I_{(3,4)}$ equals to four, the sample magnetizes completely in-plane. From the easy axis of magnetization, it can be expected that the angle θ_S for γ' -Fe₄N(001) and γ '-Fe₄N(111) would be 0° and 35.26°, respectively.

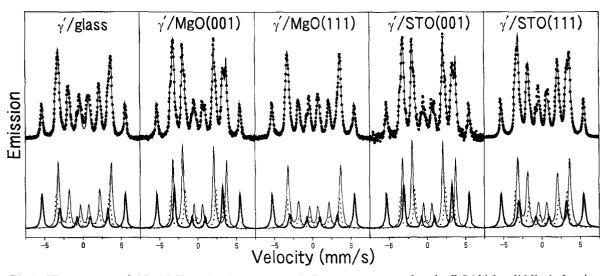


Fig.4: CEMS spectra of γ' -Fe₄N films. The lower spectra indicate Fe atoms at edge site FeI (thick solid line), fcc sites FeIIA (thin solid line) and FeIIB (dotted line).

sub doublets in	OPario ab	cona.			
	$I_{(2,5)}/I_{(3,4)}$		$\theta_{\rm S}$ (deg.)		
	(001)	(111)	(001)	(111)	
γ'/glass	1.70		39.40		
γ'/MgO	3.59	1.22	13.41	46.80	
γ'/STO	3.74	2.34	10.53	30.72	
Expected	4.00	2.00	0.00	35.26	

Table V: Angles of magnetization of γ '-Fe₄N films to the substrate surface derived from relative intensity of sub-doublets in CEMS spectra.

The values of θ_s for each film derived from Fig.4 are shown in table V. The relative intensity of the (2,5) sub-doublet of γ' /glass is about two, so the film magnetizes at random. For the (001)-oriented films and γ' /STO(001), θ_s are larger than expected value. In contrast, γ' /STO(111) has smaller θ_s than expected, as is reasonable because films tend to magnetize in in-plane direction due to the effect of demagnetizing field [8].

Fig. 5 shows the relation between rates of lattice expansion and angles θ_S for the films. If the rate is larger than a certain value, the effect of lattice expansion overcomes that of demagnetizing field and the angle of magnetization becomes larger. So we concluded that the direction of magnetization of the films except for $\gamma'/STO(111)$ is due to the large lattice expansion in the film.

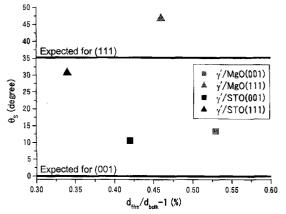


Fig.5: Relationship between lattice expansion of the films and angles of magnetization.

4. CONCLUSION

Highly oriented γ '-Fe₄N films were successfully grown on MgO(001), MgO(111), SrTiO₃(001) and SrTiO₃(111) substrates. The interplanar spacings of γ '-Fe₄N films on MgO and SrTiO₃ are enlarged in order to reduce the lattice mismatch between the films and substrates. The behavior of magnetization in the films observed by VSM and CEMS depends on the structure and crystalline orientation of the films. The (111)-oriented films show an out-of-plane magnetic anisotropy, and the films on MgO(001), MgO(111) and SrTiO₃(001) magnetize at larger angle to the substrates than expected due to the large lattice expansion.

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