# Intercalation in thermoelectric layered cobaltites

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 $Bi_2(Sr, Ca)_2Co_2O_9$  (BC-222) single crystalline whiskers intercalated  $I_2$  and HgBr<sub>2</sub> have been prepared. The *c*-cell parameter of the single crystals are elongated 3.6Å and 6.2Å by the  $I_2$  and HgBr<sub>2</sub> intercalation, respectively. The  $I_2$  and HgBr<sub>2</sub> present between two Bi-O layers. While Seebeck coefficient of the intercalated samples is lower than no-intercalated ones because of increase in hole concentration, electrical resistivity is increased by the intercalation. Moreover, the  $I_2$  intercalation reduces thermal conductivity. Key words: Thermoelectric material, Oxide, Intercalation, 2D structure

#### 1. INTRODUCTION

One form of energy expected to develop in the next few decades is thermoelectric power generated from vast amounts of waste heat emitted by automobiles, factories, and similar sources. Thermoelectric generation systems can convert heat energy into electric energy directly without using moving parts such as turbines and without producing carbon dioxide gas, radioactive substances, or other emissions. To realize thermoelectric power generation, not only higher thermoelectric figure of merit (ZT = $S^2T/\rho\kappa$ ; S: Seebeck coefficient, T: temperature,  $\rho$ : electric resistivity, and  $\kappa$ : thermal conductivity) than 1.0, but also high chemical stability and no content of harmful elements are indispensable to thermoelectric materials. For this purpose, oxide materials are the one of the strongest candidates. Especially, layered CoO<sub>2</sub>-based single crystals indicating p-type properties have enough thermoelectric figure of merit (ZT) at high temperature in air [1, 2].

Recently, improvement of thermoelectric properties has been reported by controlling nanostructure. In this strategy, it is possible only phonons are scattered but no electrons. Usually, mean free path of electrons is longer than that of electrons. In fact, huge high ZT has been obtained in Bi2Te3/Sb2Te3 superlattice thin films [3]. In this film, electrons can flow as ballistic current. On the other hand phonons are scattered at the interface between Bi2Te3 and Sb2Te3 sublattices. Although CoO<sub>2</sub>-based materials mentioned above has a superlattice structure, they shows no effect caused by the two dimensional (2D) structure obviously. It is expected to cause the 2D effect if the thickness of the rock-salt sublattice can be controlled.

In this paper,  $I_2$  and  $HgBr_2$  intercalation into the rock-salt sublattices and thermoelectric properties in intercalated  $Bi_2(Sr, Ca)_2Co_2O_9$  (BC-222) single crystals are discussed.

## 2. EXPERIMENTAL

The BC-222 single crystalline whiskers were grown by annealing glassy precursor plates possessing a cationic composition of Bi;Sr:Ca:Co = 1:1:1:2, which were obtained by quenching melt of Bi<sub>2</sub>O<sub>3</sub>, SrCo<sub>3</sub>, CaCo<sub>3</sub> and Co<sub>3</sub>O<sub>4</sub> from 1300°C in air [4]. The precursor plates were annealed at 880°C for 50h in O<sub>2</sub> gas flow. The whiskers were grown from the surface of the plates and reaped using a tweezers.

The whiskers and  $I_2$  or HgBr<sub>2</sub> were put into a glass tube and sealed in vacuum. Intercalation of  $I_2$  and HgBr<sub>2</sub> was carried out at 150°C and 400°C for 20-500h using electrical furnaces, resistively. After the heat treatment the whiskers were washed using ethanol. The *c*-cell parameter of the whiskers was measured by X-ray diffraction (XRD). Element ratio was evaluated using energy dispersive X-ray (EDX) analysis.

Seebeck coefficient (S) was measured from 90-400K using a measurement system made by MMR Technologies Inc. Electrical resistivity ( $\rho$ ) measurement was performed at 4-300K by a DC four terminals method. Thermal conductivity ( $\kappa$ ) was calculated by thermal diffusibility and specific heat measured using an AC calorimetric method and a differential scanning calorimetric method, respectively. Transmitting electron microscopic observation was carried out for the cross sections of the whiskers.

## 3. RESULT AND DISCUSSION

Figure 1 shows XRD patterns for both no-intercalated and intercalated whiskers. In the case of  $I_2$  intercalation, diffraction peaks of both original no-intercalate BC-222 and  $I_2$  intercalated BC-222 (IBC-222) phases are observed after the heat treatment for 20h at 150°C. The diffraction angles of the IBC-222 phase are lower than the BC-222 phase for the same indexed peaks. This is due to elongation of the *c*-cell parameter from 14.8Å to 18.4Å by  $I_2$  intercalation. This value is similar to that in  $Bi_2Sr_2CaCu_2O_x$  superconducting phase intercalated  $I_2$  [5].  $I_2$  penetrates between BiO-BiO double layers, which are weakly bonded by Van der Waals force (Fig. 2). Intercalation time longer than 50h, the diffraction peaks of the original BC-222 phase disappear.



Fig. 1 X-ray diffraction patterns for no-intercalated (0h) and I<sub>2</sub> intercalated whiskers for 20-100h at  $150^{\circ}$ C.



Fig. 2 Schematic crystallographic structures for no-intercalated (a) and  $I_2$  intercalated (b) BC-222 phase.

Figure 3 shows XRD patterns for the HgBr<sub>2</sub> intercalated (HBBC-222) whiskers. The *c*-cell parameter is increased by HgBr<sub>2</sub> intercalation from 14.8Å to 21.0Å. No diffraction peaks for the original BC-222 phase is observed in the whiskers intercalated

for 100h. In the case of HBBC-222 phase,  $HgBr_2$  seems to be also present between the double BiO-BiO layers.

It is clear from EDX measurement that contents of all intercalating elements increase with intercalation time. The content of I saturates at 50h, while those of Hg and Br increases up to 100h and decreases at 200h (Fig. 4).



Fig. 3 X-ray diffraction patterns for no-intercalated (0h) and HgBr<sub>2</sub> intercalated whiskers for 50 and 100h at  $400^{\circ}$ C



Fig. 4 I, Hg, and Br contents in intercalated whiskers to Co content of 2.0.

S values are decreased by both  $I_2$  and HgBr<sub>2</sub> intercalation (Fig. 5). In the case of the IBC-222 phase, this would be due to the increase in the hole concentration. It is an open question in the HBBC-222 phase. Although the hole concentration would be increased in the IBC-222 whiskers,  $\rho$  values are increased. From TEM observation, many stacking faults are formed by the  $I_2$  intercalation (Fig. 6). Electrical carriers seem to be scattered at the stacking faults. If such the stacking faults were removed,  $\rho$  could be reduced because of the increase in the hole concentration.



Fig. 5  $\rho$  and S of no-intercalated (0h) and I<sub>2</sub> intercalated (a) and HgBr<sub>2</sub> intercalated whiskers (b).



Fig. 6 A TEM image of  $I_2$  intercalated whisker for 50h.

Temperature dependence of  $\kappa$  is indicated in Fig. 7. It is clear that I<sub>2</sub> intercalation is effective to decrease  $\kappa$ . The reduction in  $\kappa$  must be caused by the stacking faults. It also might be, however, due to elongation of the rock-salt sublattice in the *c*-axis direction. *Z* values can not be improved by I<sub>2</sub> intercalation because of the decrease in *S* and increase in  $\rho$ . If I<sub>2</sub> or HgBr<sub>2</sub> are intercalated completely and the stacking faults are removed, *Z* values of the intercalated samples should become higher than the no-intercalated ones.



Fig. 7 Temperature dependence of  $\kappa$  for no-intercalated (0h) and I<sub>2</sub> intercalated whisker for 50h.

## 4. CONCLUSION

 $Bi_2(Sr, Ca)_2Co_2O_9$  (BC-222) single crystals intercalated I<sub>2</sub> and HgBr<sub>2</sub> have been prepared. The *c*-cell parameter of the single crystals is elongated 3.6Å and 6.2Å by the I<sub>2</sub> and HgBr<sub>2</sub> intercalation, respectively. Because the intercalation has been incompletely, the effect of the elongation in the *c*-axis on thermoelectric properties can not be clear. If the proper elements are intercalated without stacking faults, the enhancement of thermoelectric properties is expected by controlling nano-structure.

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